

# PHASE PORTRAITS OF THE COMPLEX ABEL POLYNOMIAL DIFFERENTIAL SYSTEMS

JAUME LLIBRE<sup>1</sup> AND CLAUDIA VALLS<sup>2</sup>

ABSTRACT. In this paper we characterize the phase portraits of the complex Abel polynomial differential equations

$$\dot{z} = (z - a)(z - b)(z - c),$$

with  $z \in \mathbb{C}$ ,  $a, b, c \in \mathbb{C}$ . We give the complete description of their phase portraits in the Poincaré disc (i.e. in the compactification of  $\mathbb{R}^2$  adding the circle  $\mathbb{S}^1$  of the infinity) modulo topological equivalence.

## 1. INTRODUCTION AND STATEMENT OF THE MAIN RESULTS

Numerous problems of applied mathematics, or in physics, chemistry, economics, ... are modeled by polynomial differential systems. Excluding linear systems the quadratic polynomial differential systems are the ones with the lowest degree of complexity, and the large bibliography on them proves their relevance, see the books [1, 12, 13] and the surveys [3, 4]. After the quadratic polynomial differential systems come the cubic ones, which also have many applications. Among the cubic polynomial differential systems we emphasize the Abel systems, see for instance the papers [2, 6, 8] where the Abel systems are applied to modelize problems from Ecology, control theory for electrical circuits and cosmology, respectively.

In this paper we characterize the phase portraits of the complex Abel differential equations

$$(1) \quad \dot{z} = (z - a)(z - b)(z - c),$$

with  $z \in \mathbb{C}$ ,  $a, b, c \in \mathbb{C}$  and the dot means derivative with respect to the time  $t \in \mathbb{R}$ . We write  $z = x + iy$ ,  $a = a_1 + ia_2$ ,  $b = b_1 + ib_2$ ,  $c = c_1 + ic_2$ , with  $x, y \in \mathbb{R}$  and  $a_i, b_i, c_i \in \mathbb{R}$  for  $i = 1, 2$ . The complex differential equation (1) becomes the real differential system

$$(2) \quad \begin{aligned} \dot{x} &= -a_1b_1c_1 + a_2b_2c_1 + a_2b_1c_2 + a_1b_2c_2 + (a_1b_1 - a_2b_2 + a_1c_1 + b_1c_1 \\ &\quad - a_2c_2 - b_2c_2)x - (a_2b_1 + a_1b_2 + a_2c_1 + b_2c_1 + a_1c_2 + b_1c_2)y \\ &\quad - (a_1 + b_1 + c_1)x^2 + 2(a_2 + b_2 + c_2)xy + (a_1 + b_1 + c_1)y^2 + x^3 - 3xy^2, \\ \dot{y} &= -a_2b_1c_1 - a_1b_2c_1 - a_1b_1c_2 + a_2b_2c_2 + (a_2b_1 + a_1b_2 + a_2c_1 + b_2c_1 \\ &\quad + a_1c_2 + b_1c_2)x + (a_1b_1 - a_2b_2 + a_1c_1 + b_1c_1 - a_2c_2 - b_2c_2)y \\ &\quad - (a_2 + b_2 + c_2)x^2 - 2(a_1 + b_1 + c_1)xy + (a_2 + b_2 + c_2)y^2 + 3x^2y - y^3. \end{aligned}$$

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2010 *Mathematics Subject Classification.* Primary 34A05. Secondary 34C05, 37C10.  
*Key words and phrases.* complex Abel system, Poincaré compactification, dynamics at infinity.

The objective of this work is to classify the phase portraits of the Abel polynomial differential systems (2) in the Poincaré disc modulo topological equivalence. As any polynomial differential system, system (2) can be extended to an analytic system on a closed disc  $\mathbb{D}$  of radius one, whose interior is diffeomorphic to  $\mathbb{R}^2$  and its boundary, the circle  $\mathbb{S}^1$ , plays the role of the infinity. This closed disc is denoted by  $\mathbb{D}^2$  and called the *Poincaré disc*, because the technique for doing such an extension is the *Poincaré compactification* for a polynomial differential system in  $\mathbb{R}^2$ , which is described in details in Chapter 5 of [5], see also subsection 2.1. In this paper we shall use the notation of that chapter. By using this compactification technique the dynamics of system (2) in a neighborhood of the infinity can be studied.

See subsection 2.2 for the definition of equivalent topological phase portraits, and for seeing that it is sufficient to draw the separatrices in the Poincaré disc and an orbit in every canonical region for characterizing a phase portrait in the Poincaré disc.

**Theorem 1.** *There are only six different topologically phase portraits of the Abel differential system (2) in the Poincaré disc, see Figure 1.*

In Figure 1 the phase portrait

- (a) is realized when  $a_1 = a_2 = -1$ ,  $b_1 = b_2 = 0$ , and  $c_1 = c_2 = -1/2$ ;
- (b) is realized when  $a_1 = a_2 = 0$ ,  $b_1 = -3/2$ ,  $b_2 = -3$ ,  $c_1 = -2$ ,  $c_2 = -1$ ;
- (c) is realized when  $a_1 = a_2 = b_1 = 0$ ,  $b_2 = 1$ ,  $c_1 = 0$ ,  $c_2 = 2$ ;
- (d) is realized when  $a_1 = a_2 = 0$ ,  $b_1 = 2$ ,  $b_2 = -1$ ,  $c_1 = -2$ ,  $c_2 = -1$ ;
- (e) is realized when  $a_1 = a_2 = 0$ ,  $b_1 = -8$ ,  $b_2 = -6$ ,  $c_1 = -3$ ,  $c_2 = -1$ ;
- (f) is realized when  $a_1 = -1$ ,  $a_2 = 0$ ,  $b_1 = -1$ ,  $b_2 = 0$ ,  $c_1 = 0$ ,  $c_2 = 1$ ;
- (g) is realized when  $a_1 = a_2 = b_1 = b_2 = 0$ ,  $c_1 = -1$ ,  $c_2 = -1/2$ ;
- (h) is realized when  $a_1 = a_2 = b_1 = b_2 = c_1 = 0$ ,  $c_2 = -1$ ; and
- (i) is realized when  $a_1 = a_2 = b_1 = b_2 = c_1 = c_2 = 0$ .

## 2. PRELIMINARIES

**2.1. Poincaré compactification.** For a complete description of the Poincaré compactification method we refer to chapter 5 of [5]. In what follows we remember some formulas.

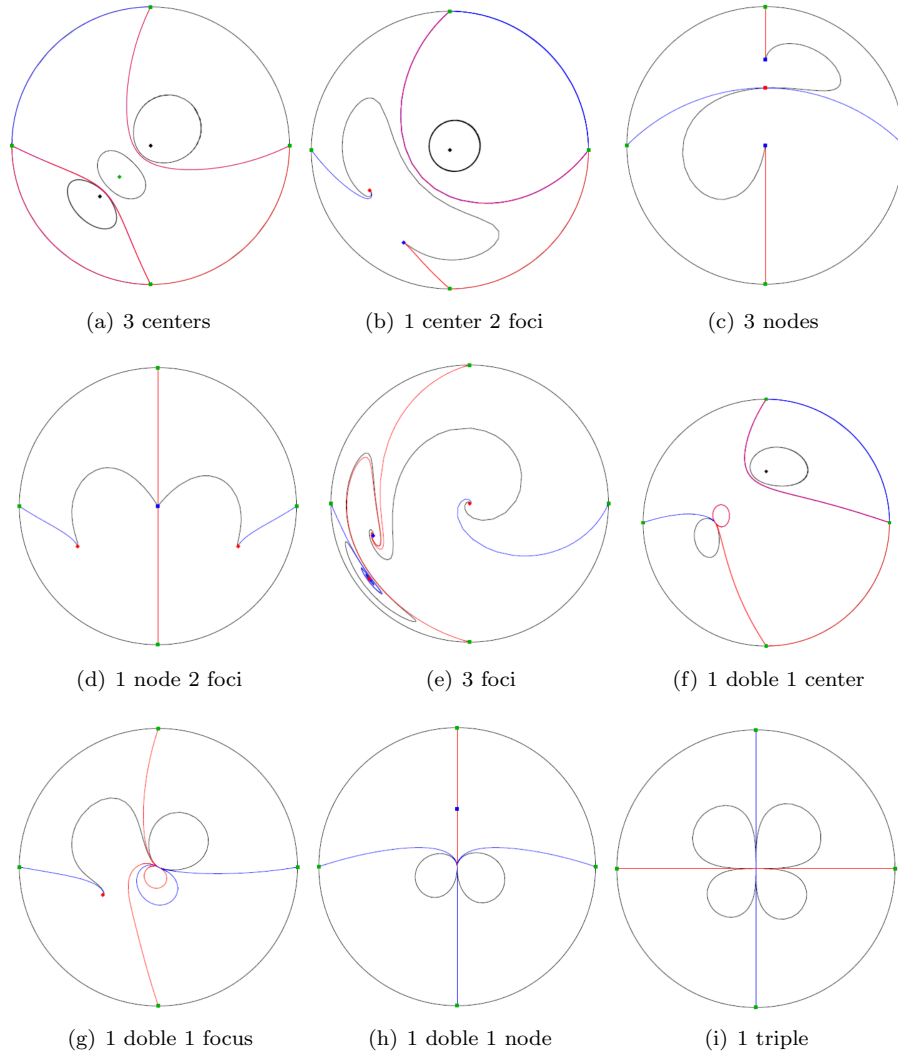
Consider a polynomial differential system in  $\mathbb{R}^2$  with degree 3.

$$\dot{x} = P(x, y), \quad \dot{y} = Q(x, y)$$

or equivalently its associated polynomial vector field  $X = (P, Q)$ . As we said before, any polynomial differential system can be extended to an analytic differential system on a closed disc of radius one centered at their origin of coordinates, whose interior is diffeomorphic to  $\mathbb{R}^2$  and its boundary, the circle  $\mathbb{S}^1$ , plays the role of the infinity.

In order to study the compactified vector field we consider four open charts covering the disc  $\mathbb{D}$ :

$$\begin{aligned} \phi_1 : \mathbb{R}^2 &\longrightarrow U_1, & \phi_1(x, y) &= (1/v, u/v), \\ \phi_2 : \mathbb{R}^2 &\longrightarrow U_2, & \phi_2(x, y) &= (u/v, 1/v) \end{aligned}$$



**Figure 1.** The nine topological phase portraits in the Poincaré disc of system (2) if we distinguish between nodes and foci, without this distinction the phase portraits (c), (d) and (e) are topologically equivalent, and also the phase portraits (g) and (h) are topologically equivalent. So there are only six different topological phase portraits.

and

$$\psi_k : \mathbb{R}^2 \longrightarrow V_k, \quad \psi_k(x, y) = -\phi_k(x, y), \quad k = 1, 2$$

with

$$\begin{aligned} U_1 &= \{(u, v) \in \mathbb{D} : u^2 + v^2 \leq 1 \text{ and } u > 0\}, \\ U_2 &= \{(u, v) \in \mathbb{D} : u^2 + v^2 \leq 1 \text{ and } v > 0\}, \\ V_1 &= \{(u, v) \in \mathbb{D} : u^2 + v^2 \leq 1 \text{ and } u < 0\}, \\ V_2 &= \{(u, v) \in \mathbb{D} : u^2 + v^2 \leq 1 \text{ and } v < 0\}. \end{aligned}$$

The Poincaré compactification is denoted by  $p(X)$ . The expression of  $p(X)$  in the chart  $U_1$  is

$$\dot{u} = v^3(-uP + Q), \quad \dot{v} = -v^4P,$$

where  $P$  and  $Q$  are evaluated at  $(1/v, u/v)$ .

The expression of  $p(X)$  in the chart  $U_2$  is

$$\dot{u} = v^3(P - uQ), \quad \dot{v} = -v^4Q,$$

where  $P$  and  $Q$  are evaluated at  $(u/v, 1/v)$ . Moreover in all these local charts the points  $(u, v)$  of the infinity have its coordinate  $v = 0$ .

The expression for the extend differential system in the local chart  $V_i$ ,  $i = 1, 2$  is the same as in  $U_i$ .

**2.2. Topological equivalence of two polynomial vector fields.** Let  $X_1$  and  $X_2$  be two polynomial vector fields on  $\mathbb{R}^2$ . We say that they are *topologically equivalent* if there exists a homeomorphism on the Poincaré disc  $\mathbb{D}$  which preserves the infinity  $\mathbb{S}^1$  and sends the orbits of  $\pi(p(X_1))$  to orbits of  $\pi(p(X_2))$ , preserving or reversing the orientation of all the orbits.

A *separatrix* of the Poincaré compactification  $\pi(p(X))$  is one of following orbits: all the orbits at the infinity  $\mathbb{S}^1$ , the finite singular points, the limit cycles, and the two orbits at the boundary of a hyperbolic sector at a finite or an infinite singular point, see for more details on the separatrices [9, 10].

The set of all separatrices of  $\pi(p(X))$ , which we denote by  $\Sigma_X$ , is a closed set (see [10]).

A *canonical region* of  $\pi(p(X))$  is an open connected component of  $\mathbb{D} \setminus \Sigma_X$ . The union of the set  $\Sigma_X$  with an orbit of each canonical region form the *separatrix configuration* of  $\pi(p(X))$  and is denoted by  $\Sigma'_X$ . We denote the number of separatrices of a phase portrait in the Poincaré disc by  $S$ , and its number of canonical regions by  $R$ .

Two separatrix configurations  $\Sigma'_{X_1}$  and  $\Sigma'_{X_2}$  are *topologically equivalent* if there is a homeomorphism  $h : \mathbb{D} \rightarrow \mathbb{D}$  such that  $h(\Sigma'_{X_1}) = \Sigma'_{X_2}$ .

According to the following theorem which was proved by Markus [9], Neumann [10] and Peixoto [11], it is sufficient to investigate the separatrix configuration of a polynomial differential system, for determining its global phase portrait.

**Theorem 2.** *Two Poincaré compactified polynomial vector fields  $\pi(p(X_1))$  and  $\pi(p(X_2))$  with finitely many separatrices are topologically equivalent if and only if their separatrix configurations  $\Sigma'_{X_1}$  and  $\Sigma'_{X_2}$  are topologically equivalent.*

### 3. INFINITE SINGULAR POINTS

The following lemma summarizes the information at the infinite singular points.

**Lemma 3.** *There are two pairs of infinite singular points which are saddles located at the origins of the local charts  $U_1$ ,  $U_2$ ,  $V_1$  and  $V_2$ .*

*Proof.* First we determine the local phase portrait of the infinite singular points in the local chart  $U_1$ . The expression of system (2) in this chart is

$$\begin{aligned}\dot{u} &= 2u - (a_2 + b_2 + c_2)v - (a_1 + b_1 + c_1)uv + (a_2b_1 + a_1b_2 + a_2c_1 + b_2c_1 \\ &\quad + a_1c_2 + b_1c_2)v^2 + 2u^3 - (a_2 + b_2 + c_2)u^2v - (a_2b_1c_1 + a_1b_2c_1 + a_1b_1c_2 \\ &\quad - a_2b_2c_2)v^3 - (a_1 + b_1 + c_1)u^3v + (a_2b_1 + a_1b_2 + a_2c_1 + b_2c_1 + a_1c_2 \\ &\quad + b_1c_2)u^2v^2 + (a_1b_1c_1 - a_2b_2c_1 - a_2b_1c_2 - a_1b_2c_2)uv^3, \\ \dot{v} &= -v + (a_1 + b_1 + c_1)v^2 + 3u^2v - 2(a_2 + b_2 + c_2)uv^2 - (a_1b_1 - a_2b_2 \\ &\quad + a_1c_1 + b_1c_1 - a_2c_2 - b_2c_2)v^3 - (a_1 + b_1 + c_1)u^2v^2 + (a_2b_1 + a_1b_2 \\ &\quad + a_2c_1 + b_2c_1 + a_1c_2 + b_1c_2)uv^3 + (a_1b_1c_1 - a_2b_2c_1 - a_2b_1c_2 - a_1b_2c_2)v^4.\end{aligned}$$

The singular points  $(u, 0)$  in the local chart  $U_1$  satisfy  $2u(1+u^2) = 0$ . So the unique singular point in the local chart  $U_1$  is the origin. Computing the eigenvalues of the Jacobian matrix at the origin we obtain that they are 2 and  $-1$ . So the origin is a saddle.

Now we analyze the phase portrait in the local chart  $U_2$ , we need to study the origin of  $U_2$ , the other infinite singular points have been studied in the local chart  $U_1$ . The expression of the system in this chart is

$$\begin{aligned}\dot{u} &= -2u + (a_1 + b_1 + c_1)v + (a_2 + b_2 + c_2)uv - (a_2b_1 + a_1b_2 + a_2c_1 + b_2c_1 \\ &\quad + a_1c_2 + b_1c_2)v^2 - 2u^3 + (a_1 + b_1 + c_1)u^2v - (a_1b_1c_1 - a_2b_2c_1 - a_2b_1c_2 \\ &\quad - a_1b_2c_2)v^3 + (a_2 + b_2 + c_2)u^3v - (a_2b_1 + a_1b_2 + a_2c_1 + b_2c_1 + a_1c_2 \\ &\quad + b_1c_2)u^2v^2 + (a_2b_1c_1 + a_1b_2c_1 + a_1b_1c_2 - a_2b_2c_2)uv^3 \\ \dot{v} &= v - (a_2 + b_2 + c_2)v^2 - 3u^2v + 2(a_1 + b_1 + c_1)uv^2 - (a_1b_1 - a_2b_2 + a_1c_1 \\ &\quad + b_1c_1 - a_2c_2 - b_2c_2)v^3 + (a_2 + b_2 + c_2)u^2v^2 - (a_2b_1 + a_1b_2 + a_2c_1 + b_2c_1 \\ &\quad + a_1c_2 + b_1c_2)uv^3 + (a_2b_1c_1 + a_1b_2c_1 + a_1b_1c_2 - a_2b_2c_2)v^4.\end{aligned}$$

Note that the origin of the local chart  $U_2$  is a singular point whose eigenvalues of the Jacobian matrix at this point are  $-2$  and  $1$ . So the origin is a saddle. This completes the proof of the lemma.  $\square$

#### 4. FINITE SINGULAR POINTS

The finite singular points of system (2) are the real solutions of  $\dot{x} = \dot{y} = 0$ . Computing such solutions we obtain

$$(x, y) = (a_1, a_2), (x, y) = (b_1, b_2), (x, y) = (c_1, c_2).$$

Now we study the local phase portrait of each of the finite singular points.

We will use the following three propositions.

**Proposition 4.** *Let  $p$  be a singular point of system (2) such that their eigenvalues are purely imaginary. Then  $p$  is an isochronous center.*

The proof of Proposition 4 is given in [7, Corollary 2.7 (a)].

**Proposition 5.** *Let  $p$  be a singular point of system (2) of multiplicity  $k \in \{2, 3\}$ . Then the local phase portrait of  $p$  is formed by  $2(k-1)$  elliptic sectors separated by parabolic sectors.*

The proof of Proposition 5 is given in [7, Corollary 2.7 (b)].

**Proposition 6.** *System (2) has no limit cycles.*

The proof of Proposition 6 is given in [7, Theorem 1.3 (a)].

*The finite singular point  $(a_1, a_2)$*

The eigenvalues associated to the finite singular point  $(a_1, a_2)$  are  $\lambda_1 \pm i\lambda_2$  where

$$\begin{aligned}\lambda_1 &= a_1^2 - a_2^2 - a_1b_1 + a_2b_2 - a_1c_1 + b_1c_1 + a_2c_2 - b_2c_2 \\ \lambda_2 &= b_2c_1 - a_2(b_1 + c_1) + b_1c_2 - a_1(b_2 + c_2 - 2a_2).\end{aligned}$$

We have different possibilities. If  $\lambda_2 \neq 0$  then we have a focus, hyperbolic if  $\lambda_1 \neq 0$  and by Proposition 4 is a center if  $\lambda_1 = 0$ . If  $\lambda_2 = 0$  and  $\lambda_1 \neq 0$  then we have a node. If  $\lambda_1 = \lambda_2 = 0$  then the singular point is linearly zero. In this case we have that either  $b = a$  and  $c \neq a$ , or  $b = c = a$ . In view of Proposition 5 in the first case the singular point  $(a_1, a_2)$  is formed by two elliptic sectors separated by parabolic sectors and in the second case the singular point  $(a_1, a_2)$  is formed by four elliptic sectors separated by parabolic sectors.

*The finite singular point  $(b_1, b_2)$*

The eigenvalues associated to the finite singular point  $(b_1, b_2)$  are  $\lambda_3 \pm i\lambda_4$  where

$$\begin{aligned}\lambda_3 &= -a_1b_1 + b_1^2 + a_2b_2 - b_2^2 + a_1c_1 - b_1c_1 - a_2c_2 + b_2c_2 \\ \lambda_4 &= a_2b_1 + a_1b_2 - 2b_1b_2 - a_2c_1 + b_2c_1 - a_1c_2 + b_1c_2.\end{aligned}$$

We have different possibilities. If  $\lambda_4 \neq 0$  then we have a focus, hyperbolic if  $\lambda_3 \neq 0$  and by Proposition 4 is a center if  $\lambda_3 = 0$ . If  $\lambda_4 = 0$  and  $\lambda_3 \neq 0$  then we have a node. Finally, if  $\lambda_4 = 0$  and  $\lambda_3 = 0$ , then the singular point is linearly zero. We can consider that  $b \neq a$  because this case has been already studied. Then the condition  $\lambda_4 = \lambda_3 = 0$  implies  $b = c$ . Since  $c \neq a$  then in view of Proposition 5 the singular point  $(b_1, b_2)$  is formed by two elliptic sectors separated by parabolic sectors.

*The finite singular point  $(c_1, c_2)$*

The eigenvalues associated to the finite singular point  $(c_1, c_2)$  are  $\lambda_5 \pm i\lambda_6$  where

$$\begin{aligned}\lambda_5 &= a_1b_1 - a_2b_2 - a_1c_1 - b_1c_1 + c_1^2 + a_2c_2 + b_2c_2 - c_2^2, \\ \lambda_6 &= b_2c_1 + a_2(-b_1 + c_1) + b_1c_2 - 2c_1c_2 + a_1(-b_2 + c_2).\end{aligned}$$

We have different possibilities. If  $\lambda_6 \neq 0$  then we have a focus, hyperbolic if  $\lambda_5 \neq 0$  and by Proposition 5 is a center if  $\lambda_5 = 0$ . If  $\lambda_6 = 0$  and  $\lambda_5 \neq 0$  then we have a node. Finally, if  $\lambda_6 = \lambda_5 = 0$  then either  $c = a$ , or  $c = b$ , and these two cases have been studied before.

Combining the different possibilities for the finite singular points we have in principle the following cases:

three centers;

two centers and a node;

two centers and a focus;  
 one center and two nodes;  
 one center, one node and one focus;  
 one center and two foci;  
 three nodes;  
 two nodes and a focus;  
 one node and two foci;  
 three foci;  
 one point formed by two elliptic sectors separated by parabolic sectors and a center;  
 one point formed by two elliptic sectors separated by parabolic sectors and a node;  
 one point formed by two elliptic sectors separated by parabolic sectors and a focus;  
 and  
 one point formed by four parabolic sectors separated by parabolic sectors.

However the cases:

two centers and a node;  
 two centers and a focus;  
 one center and two nodes;  
 one center, one node and one focus; and  
 two nodes and a focus

are not possible. We only show that the case two centers and a node is not possible, the other non possible cases are proved in a similar way. Indeed, the singular points  $a$  and  $b$  are centers if and only if  $\lambda_1 = \lambda_3 = 0$  and  $c$  is a node if and only if  $\lambda_6 = 0$ . Then it is easy to check that the system  $\lambda_1 = \lambda_3 = \lambda_6 = 0$  has no solutions in  $a_i$ ,  $b_i$  and  $c_i$  for  $i = 1, 2$ .

In short, only nine possibilities for the finite singular points are indeed realizable, the ones of Figure 1.

## 5. GLOBAL PHASE PORTRAITS

Combining the study of the local phase portraits of the finite and infinite singular points, together with Proposition 6, which states that the Abel system (2) has no limit cycles, and combining the different possibilities for the finite singular points we conclude that the possible global phase portraits are the nine described in Figure 1.

## ACKNOWLEDGEMENTS

The first author is supported by the Ministerio de Economía, Industria y Competitividad, Agencia Estatal de Investigación grants PID2019-104658GB-I00 (FEDER), the Agència de Gestió d'Ajuts Universitaris i de Recerca grant 2017SGR1617, and the H2020 European Research Council grant MSCA-RISE-2017-777911. The second author is partially supported by FCT/Portugal through UID/MAT/04459/2019.

## REFERENCES

- [1] J.C. ARTÉS, J. LLIBRE, D. SCHLOMIUK AND N. VULPE, *Geometric configurations of singularities of planar polynomial differential systems. A global classification in the quadratic case*, to appear in Birkhäuser.

- [2] D.M. BENARDETE, V.W. NOONBURG AND B. POLLINA, *Qualitative tools for studying periodic solutions and bifurcations as applied to the periodically harvested logistic equation*, Amer. Math. Monthly **115** (2008), 202–219.
- [3] C. CHICONE AND T. JINGHUANG, *On General Properties of Quadratic Systems*, The American Mathematical Monthly, **89**, No. 3 (1982), 167–178.
- [4] W.A. COPPEL, *A Survey of Quadratic Systems*, J. Differential Equations **2** (1966), 293–304.
- [5] F. DUMORTIER, J. LLIBRE AND J.C. ARTÉS, *Qualitative Theory of Planar Differential Systems*, Springer, 2006.
- [6] E. FOSSAS, J.M. OLM AND H. SIRA-RAMÍREZ, *Iterative approximation of limit cycles for a class of Abel equations*, Phys. D **237** (2008), 3159–3164.
- [7] A. GARIJO, A. GASULL AND X. JARQUE, *Local and global phase portrait of equation  $\dot{z} = f(z)$* , Discrete Contin. Dyn. Syst. **17** (2007), 309–329.
- [8] T. HARKO AND M.K. MAK, *Relativistic dissipative cosmological models and Abel differential equation*, Comput. Math. Appl. **46** (2003), 849–853.
- [9] L. MARKUS, *Global structure of ordinary differential equations in the plane*: Trans. Amer. Math Soc. **76** (1954), 127–148.
- [10] D. A. NEUMANN, *Classification of continuous flows on 2-manifolds*, Proc. Amer. Math. Soc. **48** (1975), 73–81.
- [11] M.M. PEIXOTO, *On the classification of flows on 2-manifolds*. Academic, New York, pages 389–419, 1973. Dynamical systems (Proc. Sympos., Univ. Bahia, Salvador, 1971).
- [12] J.W. REYN, *Phase portraits of planar quadratic systems*, *Mathematics and its Applications*, **583**, Springer, 2007.
- [13] YE YANQIAN ET AL., *Theory of limit cycles*, Trans. of Mathematical Monographs **66**, Amer. Math. Soc., Providence, RI, 2 edition, 1984.

<sup>1</sup> DEPARTAMENT DE MATEMÀTIQUES, UNIVERSITAT AUTÒNOMA DE BARCELONA, 08193 BEL-LATERRA, BARCELONA, CATALONIA, SPAIN

*Email address:* jllibre@mat.uab.cat

<sup>2</sup> DEPARTAMENTO DE MATEMÁTICA, INSTITUTO SUPERIOR TÉCNICO, UNIVERSIDADE DE LISBOA, AV. ROVISCO PAIS 1049-001, LISBOA, PORTUGAL

*Email address:* cvalls@math.tecnico.ulisboa.pt